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THE FLUX AND ENERGY SPECTRUM OF PRIMARY COSMIC RAY ELECTRONS*

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Since the discovery of primary electrons in the cosmic radiation 3,4,5,6 attempts have been made to measure the flux of these particles at various energies.

The low intensity of the electron component causes great difficulties in obtaining enough events for good statistics and energy resolution and has so far prevented the determination of a reliable energy spectrum. Yet, the knowledge of the electron spectrum in the vicinity of the earth is important for several reasons. At the present time, it appears likely that the primary electrons observed near the earth are of galactic origin. A determination of their flux and spectrum makes it possible to investigate in detail their relation to the nonthermal galactic radio emission. Furthermore, the modulation of their intensity and spectrum due to solar controlled mechanisms may be different from the modulation of heavier primary particles. Parker has discussed the possibility of velocity dependent modulation, which can be tested by investigating the primary cosmic ray electrons.

The experiment which we shall discuss here gives an approximate energy spectrum of the primary electron component. It was carried out in the summer of 1964 at a period near solar activity minimum. This period is most advantageous for

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investigations of low energy galactic particles since the effects of solar modulation are approaching their minimum. Two balloon flights were made at Ft. Churchill on July 22 and 29, 1964, both of them floating under about 4g/cm² of residual atmosphere, in order to measure the flux and energy spectrum of the electron component. A schematic cross section of the instrument which was used is shown in Figure 1. Vertically incident particles trigger a counter telescope consisting of the scintillation counters T and I and a gas-Cerenkov counter C. The geometry factor of the telescope is 1.08 cm²sr. The gas counter is filled with 2.4 atm. of SF_A and has a threshold of about 8 MeV for electrons and 15 BeV for protons. The energy loss of the particles in Counter I is measured in order to discriminate between singly and multiply charged particles. After passing through Counter I, the particles penetrate a layer of 10.8 g/cm² of lead, a plastic scintillation counter S, and enter a lead glass Cerenkov Counter II (Schott glass No. SF6-FA) with a total depth of 13 radiation lengths. The amount of Cerenkov light produced in this counter is measured by a pulse height analyzer. Provided the incident particle is an electron and the photon-electron shower produced by this electron is absorbed in the lead glass, then the light output is a measure of the energy of the incident electron. The lead glass Cerenkov counter is completely surrounded by anticoincidence counters in order to discriminate against events in which one or several particles emerge from the lead glass. All events in which the anticoincidence counter A3 is fired are also recorded and high energy protons which penetrate the entire telescope are used for in-flight calibration of the Counters I and II. In order to further discriminate against high energy protons and interactions produced by protons in the lead glass, the scintillation counter S is

connected to two independent triggers which respond to singly charged minimum ionizing particles and to particles with 1.6 times minimum ionization or more, respectively. The equipment was calibrated in a monoenergetic beam of electrons* at energies ranging from 700 MeV to 4 BeV. These calibration runs served to determine: a) the pulse height in Counter II as a function of the energy of the incident electrons; b) the resolution of the lead glass Cerenkov counter at various energies; c) the resolution of Counter I; d) the fraction of events in which any of the anticoincidence counters was triggered as a function of the energy of the incident electrons; e) the probability that an incident electron produces more than one particle in Counter S as a function of energy; and f) the efficiency of the gas Cerenkov counter which was measured to be 95 percent. The resolution of Counter I is 37 percent for relativistic electrons (full width at half maximum) and the lead glass Cerenkov counter has a resolution which varies slightly with energy. We obtained 32 percent at 700 MeV to 21 percent at 1.3 BeV. Using the calibration data it is possible to show that contributions by protons to the measured electron flux are less than 5 percent at any energy. In Fig. 2 we present the electron energy spectrum under 4.1 g/cm² of residual atmosphere which we obtained from 20 hours of exposure at ceiling (both flights combined) and which is based on 330 electrons with energies exceeding 180 MeV.

In order to arrive at the spectrum of the primary electrons, the contribution of secondary electrons which are produced in the residual atmosphere above the equipment has to be estimated. We have, at the present time, based this estimate on

^{*}We wish to express our gratitude to Dr. Livingston for making a beam of the Cambridge Electron Accelerator available to us and to Drs. Fotino, Hand and Engels for help in setting up our experiment.

the following data: 1) The measured altitude dependence of the electron flux; 2) Simplified calculations based on the cross section and multiplicity of π -meson production by cosmic ray protons, and 3) Measurements of the μ -meson energy spectrum obtained at aircraft altitude. 8,9 The resulting energy spectrum of secondary electrons under 4.1 g/cm of atmosphere is shown as a dashed line in Figure 2. It should be noted that the estimate of the secondary electron flux is preliminary. We shall investigate it in greater detail in the future.

At energies exceeding about 500 MeV the secondary correction becomes small compared to the measured flux of electrons. In this region our results will be little modified by a more refined analysis. We, therefore, restrict this discussion to primary electrons with energies above 500 MeV. In Figure 3 the energy spectrum of primary electrons in the range from 500 MeV to 3 BeV is shown after applying corrections for secondaries and for energy loss (ionization and bremsstrahlung). We may represent this spectrum as a power law of the form

$$\frac{dJ}{dE} = 11 \times E^{-1.6} \text{ (m}^2 \text{ sec sr Bev)}^{-1}$$

where E is measured in BeV. The exponent has an error of about ± 0.5. This spectrum falls inside the limits set by earlier determinations 1,2,5 and, if extrapolated beyond 4.5 BeV, is compatible with the results of Agrenier et al³. An extrapolation to lower energies falls closely on the data which Cline et al⁶ obtained below 10 MeV. Although it is interesting to note this point, it does not lead to any conclusions at this time. In view of the large error limits in the earlier data , it is not surprising that solar modulation effects, even if present, cannot be noticed in the comparison with the results of this paper.

In the energy interval which we cover in this experiment, the exponent of the power law is appreciably smaller than would be expected on the basis of a collision origin of the electron components in the galaxy 10.11. It is lower, although not incompatible, with the exponent deduced from observation of the frequency spectrum of the nonthermal galactic radio emission.

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Mr. R. Eckstrom provided the computer program for the analysis of the data.

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Figure Captions

- Figure 1. Schematic Cross-section of the Counter Telescope
- Figure 2. The energy spectrum of electrons as measured under 4.1 g/cm² of residual atmosphere

Experimental Points

- ---- Estimated flux and energy spectrum of secondary electrons.
- Figure 3. The flux and energy spectrum of primary electrons between 500 MeV and 3.5 BeV.

(corrected for secondary electrons and energy loss in the atmosphere)





